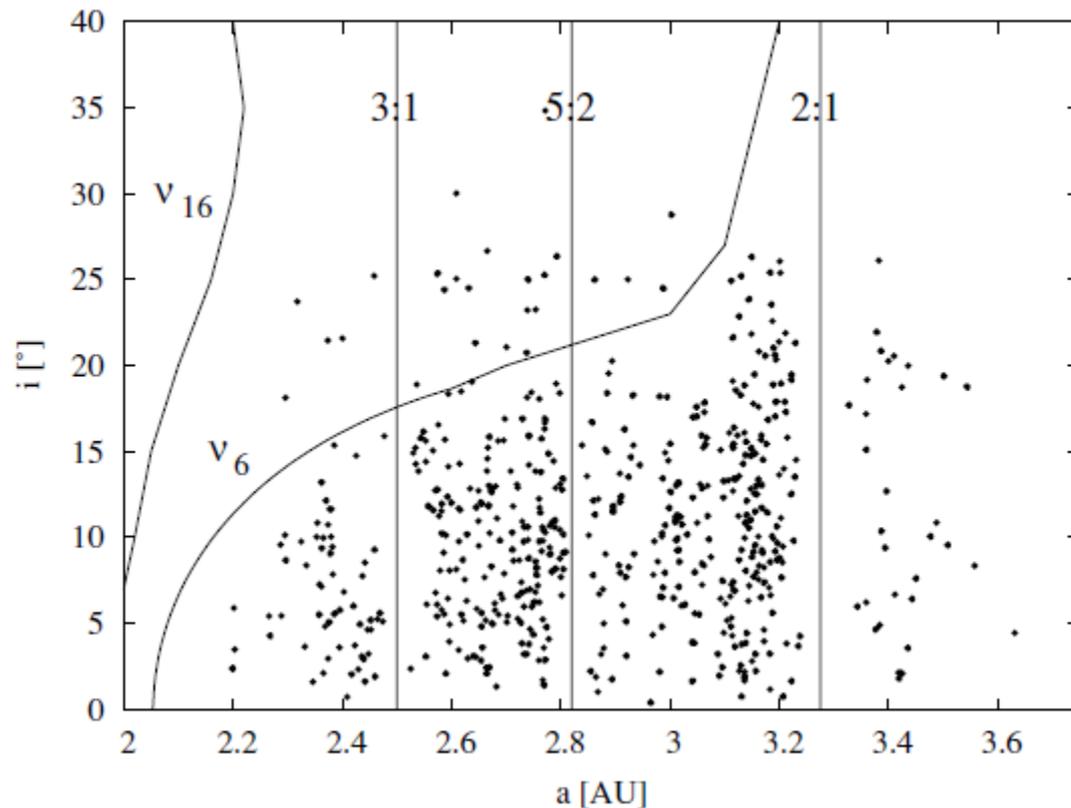


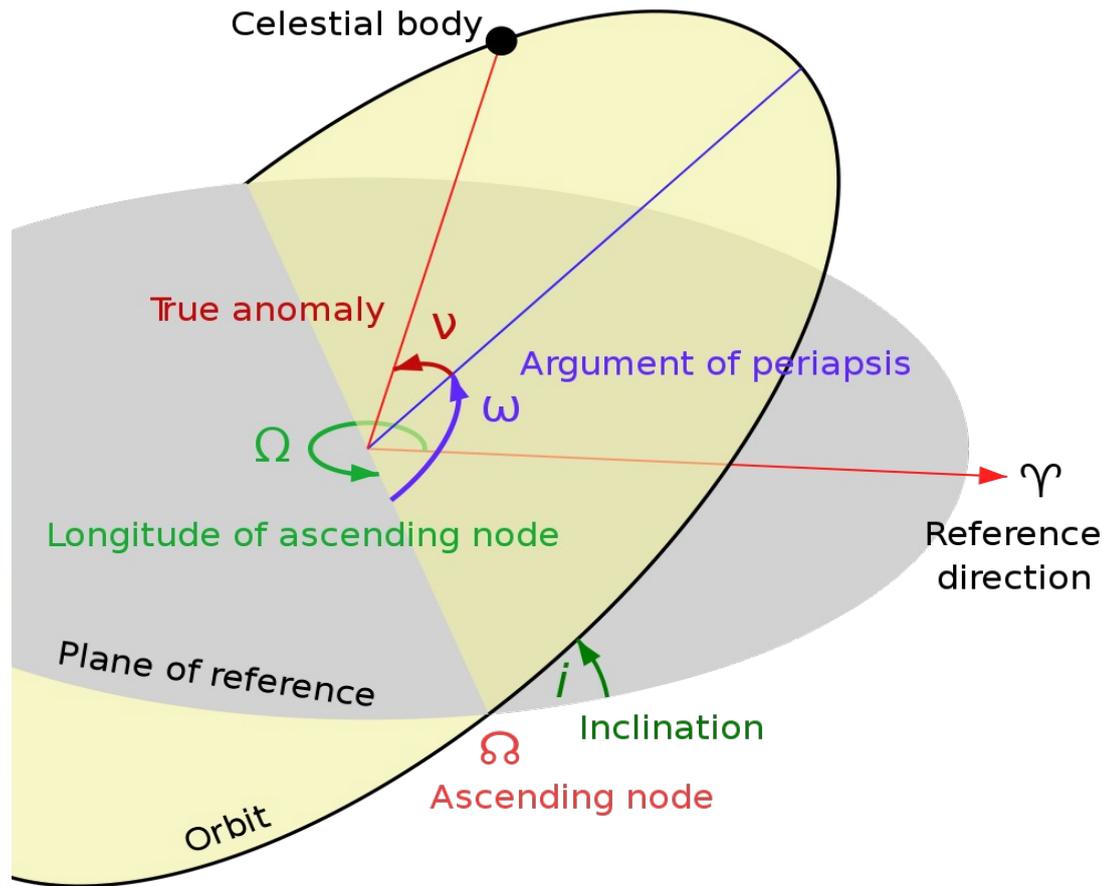
Origin of the orbital distribution of the main-belt asteroids

Kleomenis Tsiganis

*Section of Astrophysics, Astronomy & Mechanics
Department of Physics, Aristotle University of Thessaloniki, Greece*



Overview



- Orbits characterized by the elliptic orbital elements:

$$[a, e, i, \Omega, \omega, v(t)]$$

→ only **constant** in the Keplerian approximation

- The orbital distribution of *main-belt asteroids* ($D > 50$ km) contains much of the information related to the **early dynamical evolution** of the solar system

Orbital distribution in (a, e, i)

Boundaries (and gaps) are related to the main **resonances**

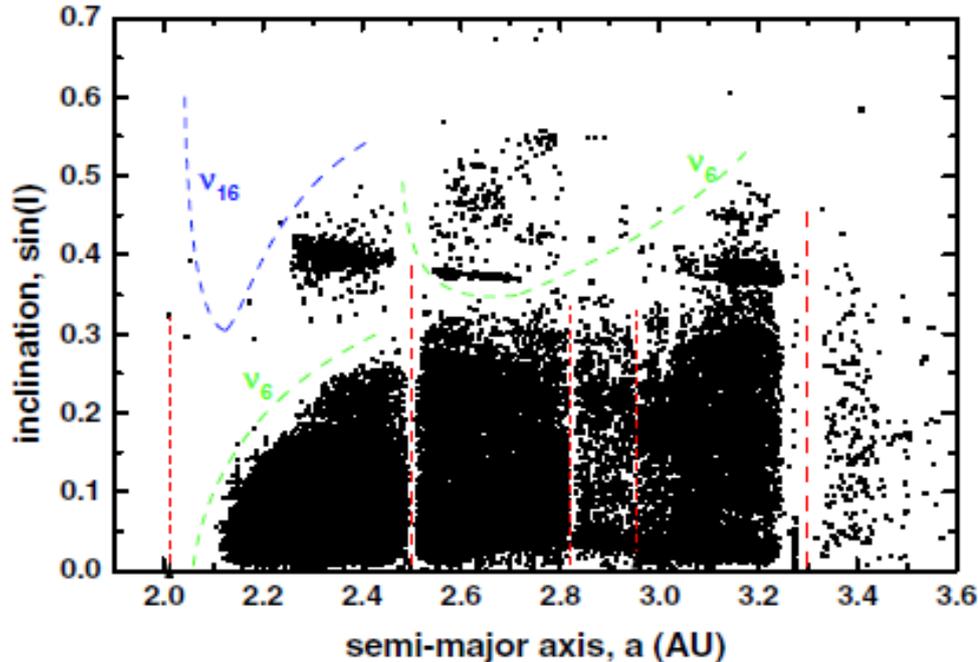
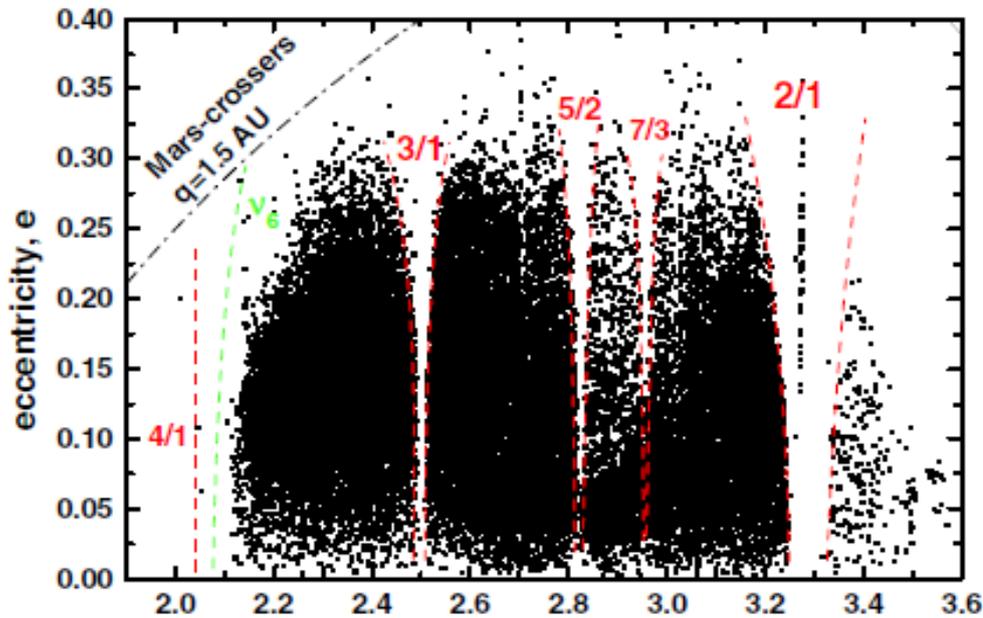
- **Kirkwood gaps** → mean motion resonances:

$$n/n_j = k/q$$

- Outer boundaries carved by **secular resonances**:

$$g = \frac{d\omega}{dt} = \langle \omega_{J,S} \rangle \longrightarrow \nu_{5,6}$$

$$s = \frac{d\Omega}{dt} = \langle \Omega_S \rangle \longrightarrow \nu_{16}$$

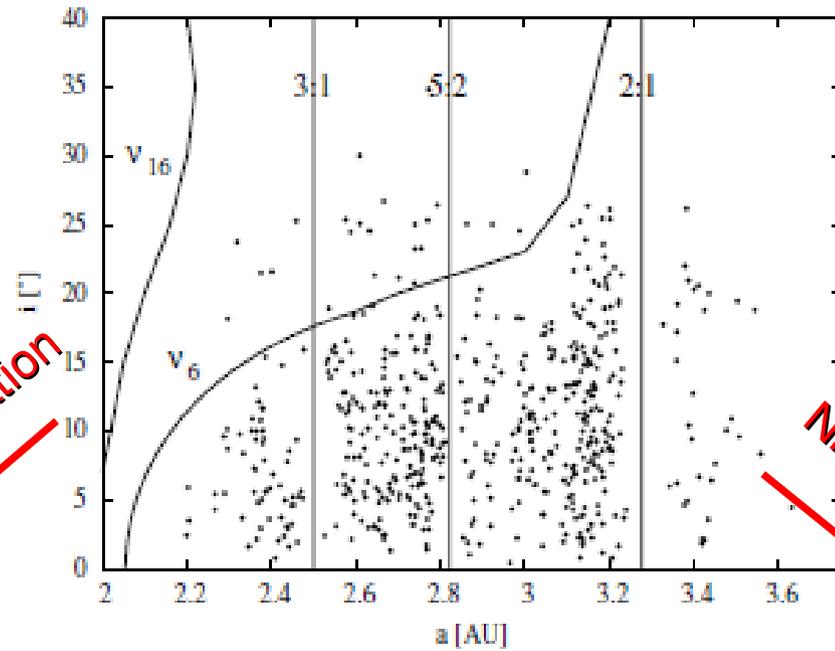


What do we need to explain?

- **Fact 1:** a “flat disk” **cannot evolve** into what we observe today
→ **need for planet migration!**
- **Fact 2:** “Smooth” migration **can explain** the (a,e) distribution
but not the (a,i) distribution → **need for faster migration**
 - Excitation is due to **Secular Resonance Sweeping:**
 $g=g(a/a_j)$, $s=s(a/a_j)$ and a_j changes with time, the main secular resonances **“sweep”** through the belt, exciting asteroid orbits (if the time-scale is of ~ 10 My)
- **Fact 3:** The “Nice model” suggests that the **(a,i) distribution is primordial**, since planetary orbits change so fast ($t < 1$ My) that asteroids do not “see” the SRs
 - **We need to find a pre-migration process to do this!!**

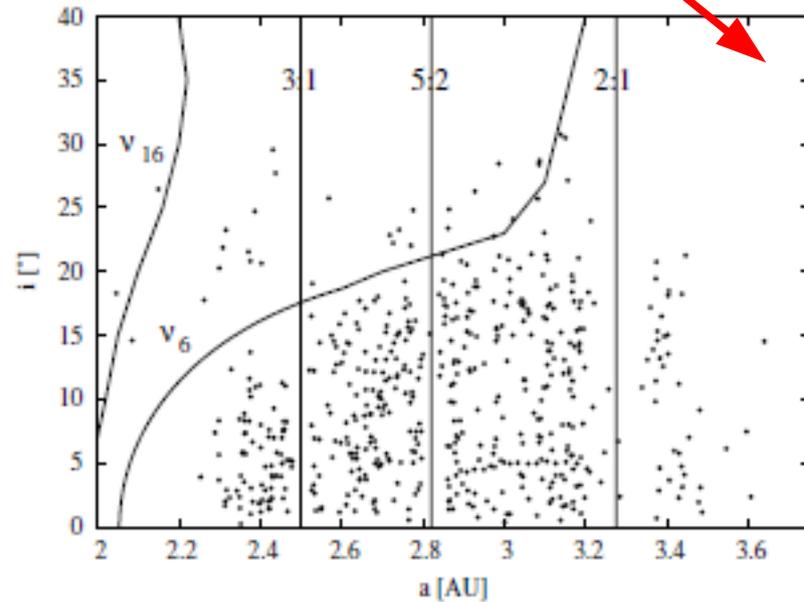
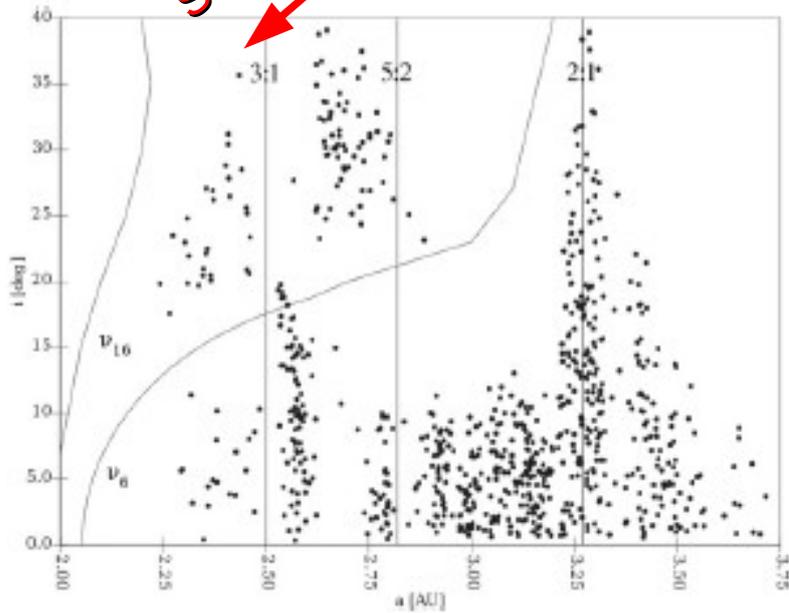
(a,i) distribution

Main-belt asteroids with $D > 50$ km



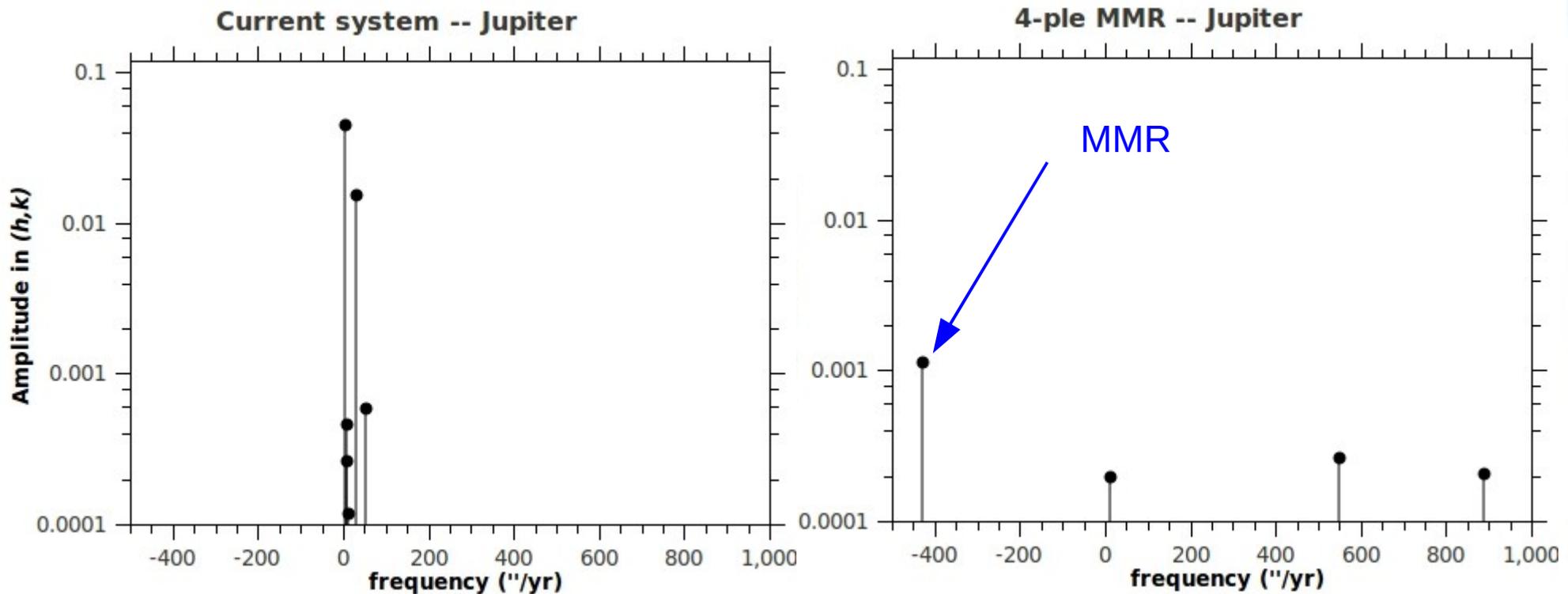
Smooth (slow) migration

Nice-model (fast)



Pre-migration planetary history

- Hydrodynamical simulations suggest that, during the **gas-rich** phase of the solar system, the planets migrate towards a **quadruple mean motion resonance** (3:2, 3:2, 4:3 /...)
- This configuration has a completely different set of **fundamental secular frequencies** w.r.t. the current system

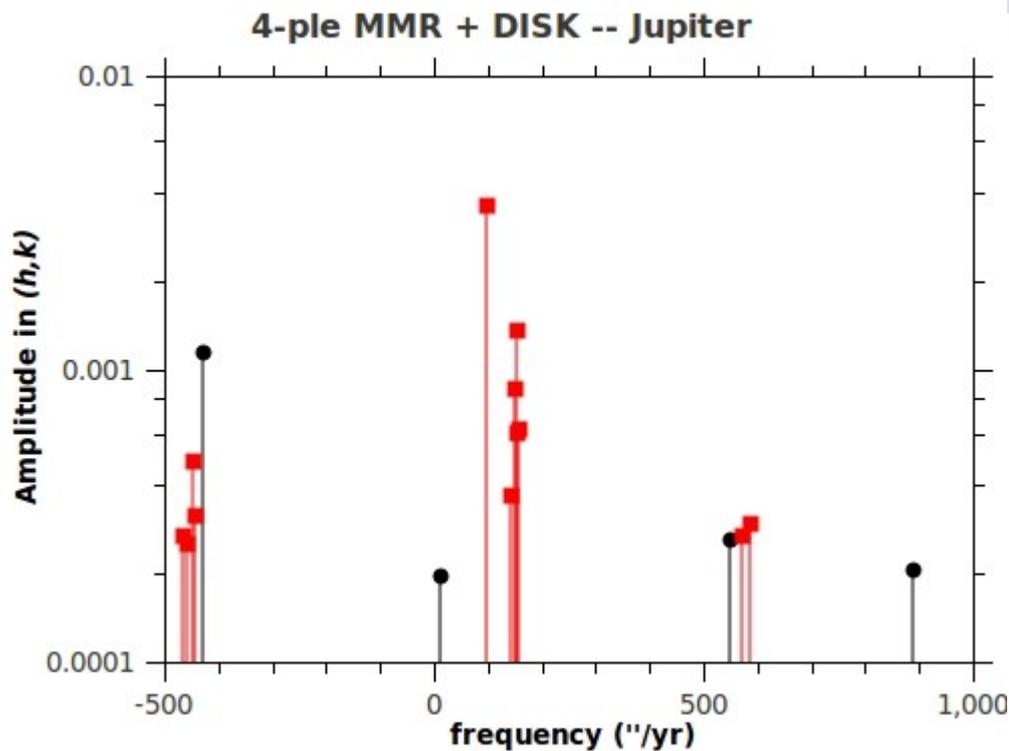


Adding a $\sim 5 M_J$ gas disk

- The pericentre frequency of an asteroid (i.e. test-particle) inside a massive disk becomes **negative** (regression)

$$n = \frac{1}{a} \left(\frac{\partial \Phi}{\partial r} \right)$$
$$\kappa = \frac{3}{a} \left(\frac{\partial \Phi}{\partial r} \right) + \left(\frac{\partial^2 \Phi}{\partial r^2} \right) \quad \longrightarrow \quad g = n - \kappa < 0$$

- assume a simple exponential evaporation model for the disk, with a characteristic e-folding time of $\tau = 0.5 - 10$ My
 - planets must go **from precession** (due to the massive disk) **to regression** (due to the resonance), as the disk evaporates
 - asteroids must go **from regression** to (slow) positive **precession** (only planets)
- frequencies on **“collision course”** → **secular resonances (SR) sweeping**



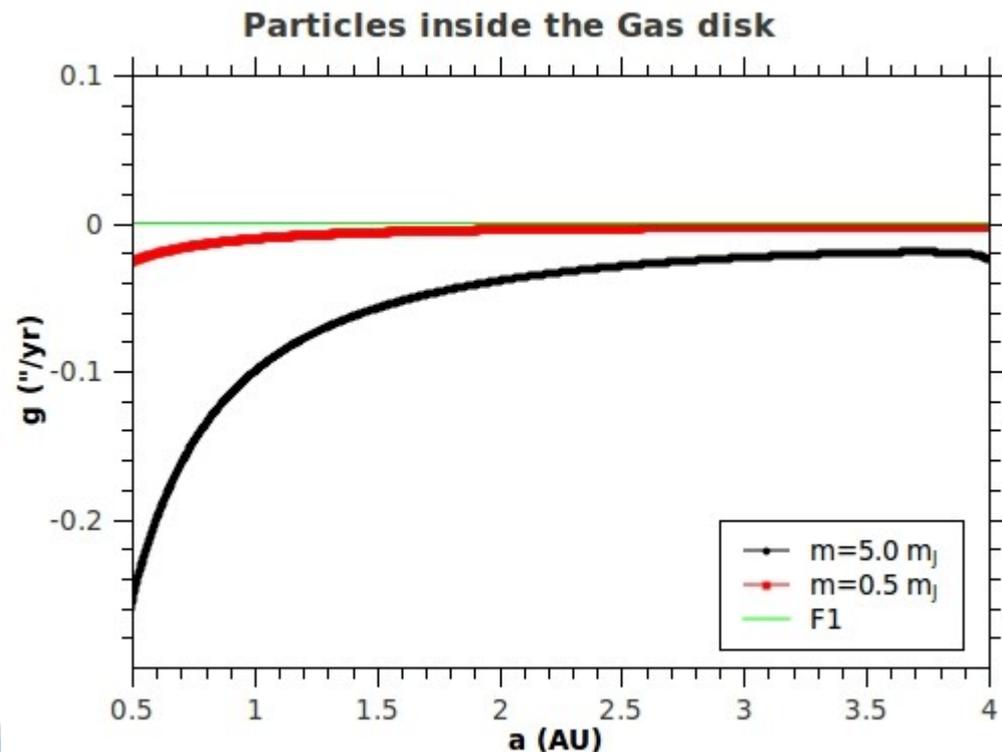
Adding a $\sim 5 M_J$ gas disk

MMSN suggests $\sim 5 M_J$ within ~ 4 AU

The spectra of the planetary orbits change as **power shifts** to $g_i > 0$

The asteroids **regress** ($g < 0$), until the mass of the disk becomes very small

→ forcing by the planets **dominates** orbital precession



The experiments

- We use the SWIFT **symplectic integrator**, adding the radial and vertical components of the **force exerted by a thin disk**:

$$\rho(R, z) = \rho_0 R^{-a} \exp[-(z/H)^b] \quad , \quad H \sim 0.01$$

- We included **tidal eccentricity damping** and **gas drag** → orbits of $D \sim 100$ km asteroids **are not** much affected
- **Two** different disk models gave **very similar** results

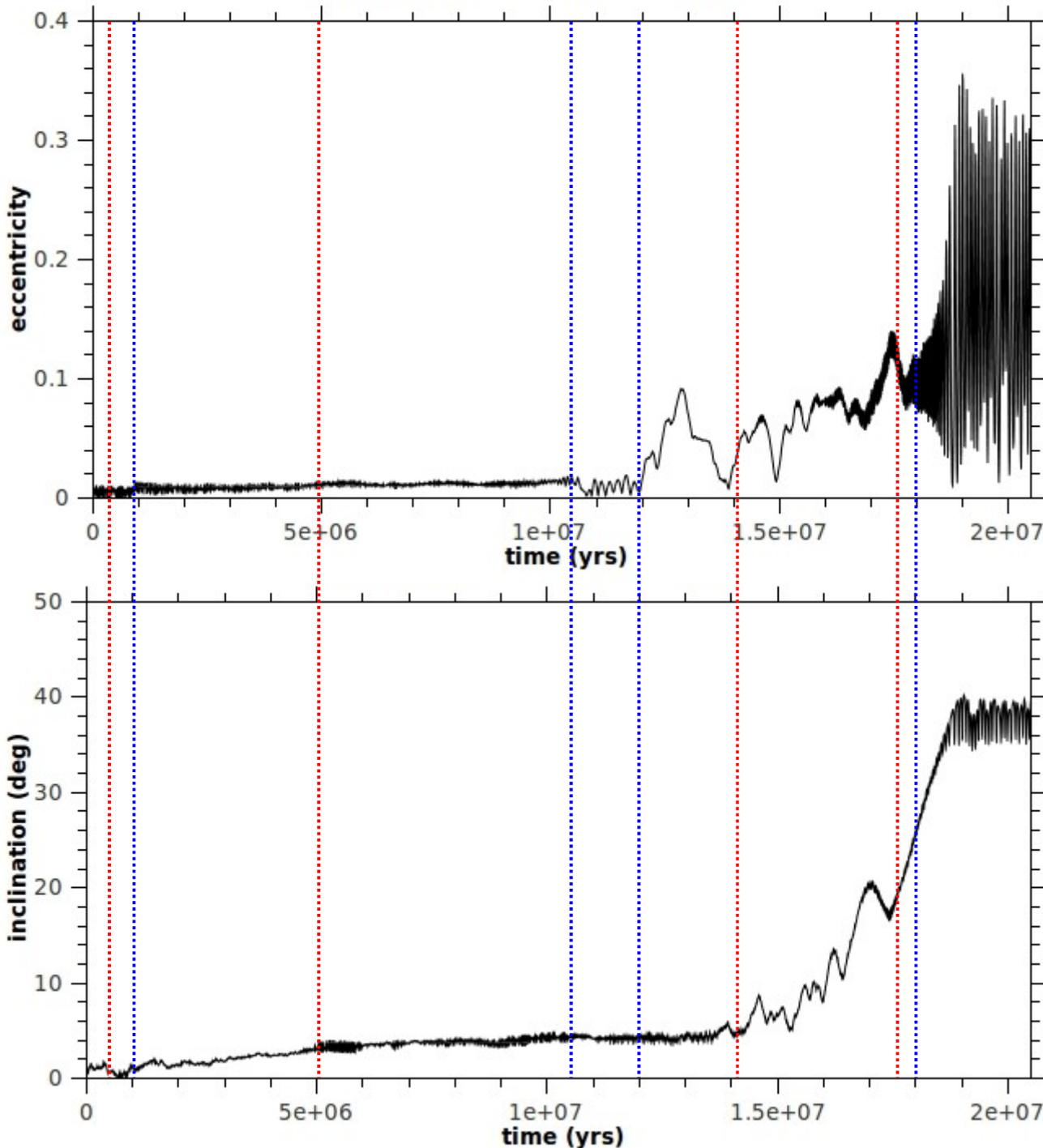
a) a **razor-thin disk** ($b=2$, $a=3/2$) for which F_z is given by:

$$4\pi G \rho(R, z) = \nabla^2 \Phi \approx \frac{\partial^2 \Phi}{\partial z^2}$$

b) a thin **exponential disk** ($b=1$) for which a full 2-D computation of $\Phi(R, z)$ was done

- **Critical** quantity is the evaporation time-scale:
 - if **too fast**, excitation but **no mixing** in e, i

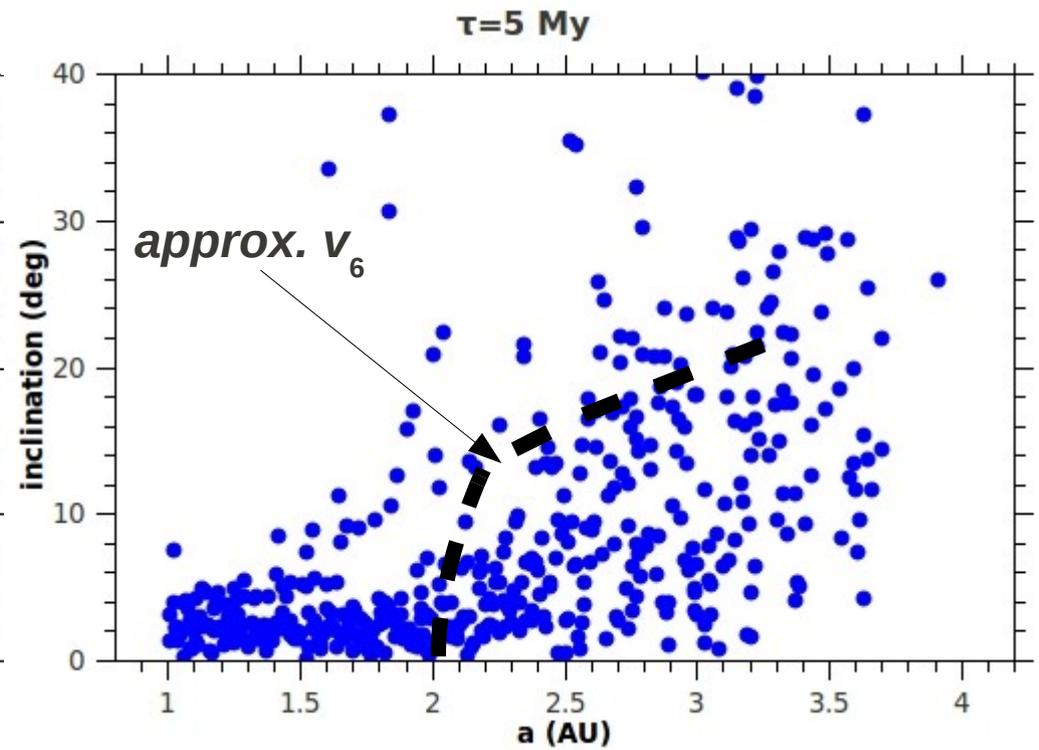
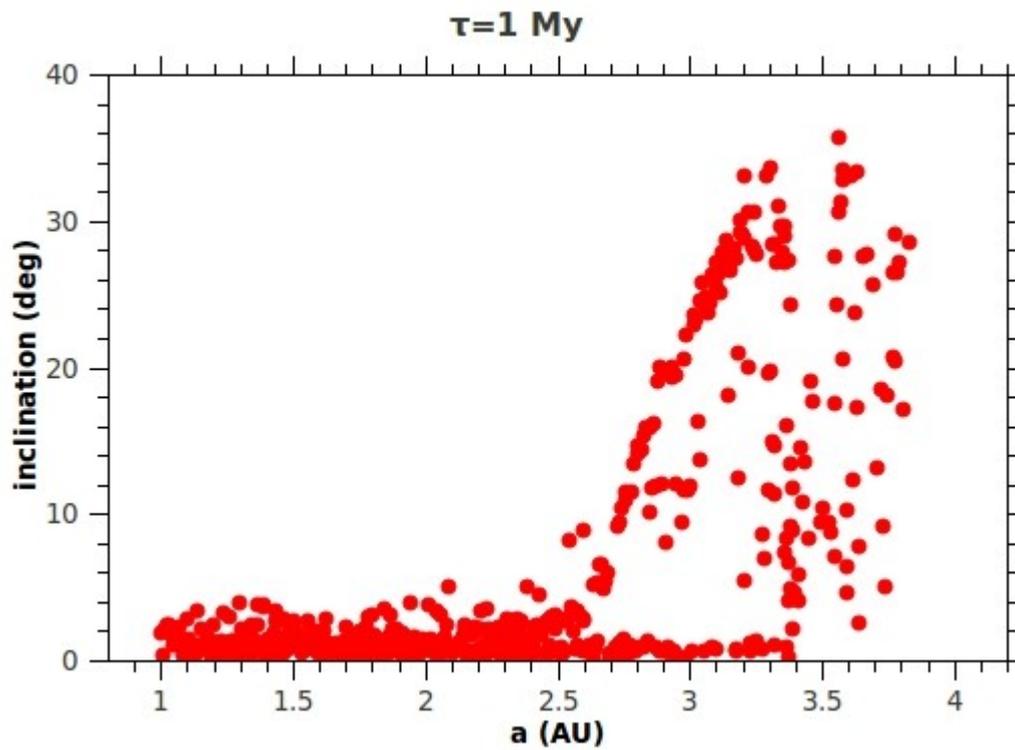
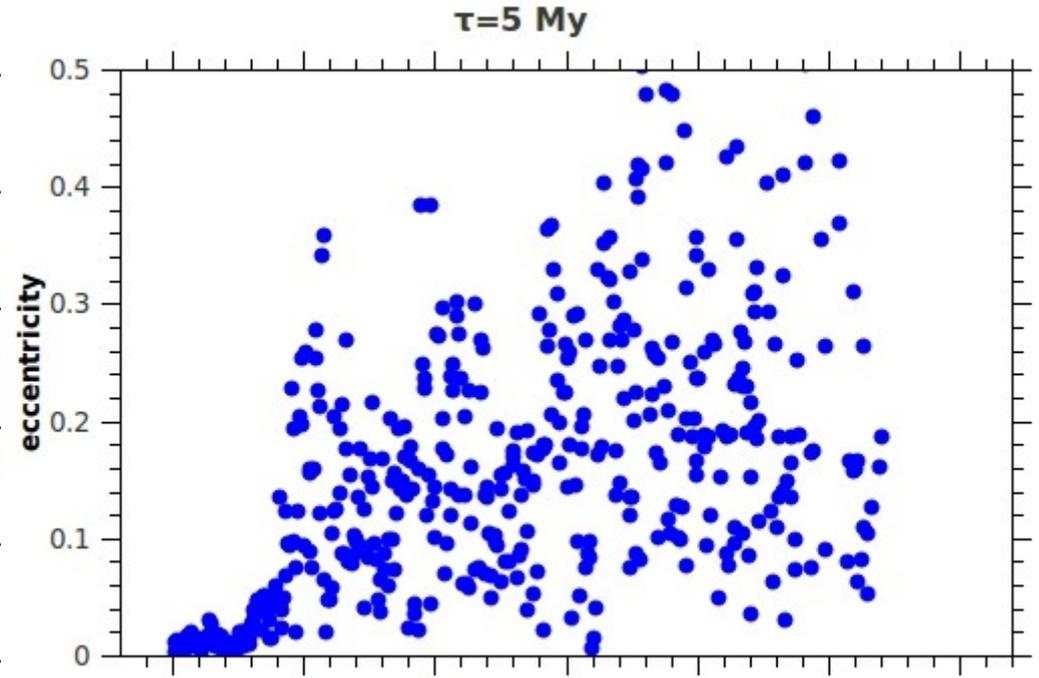
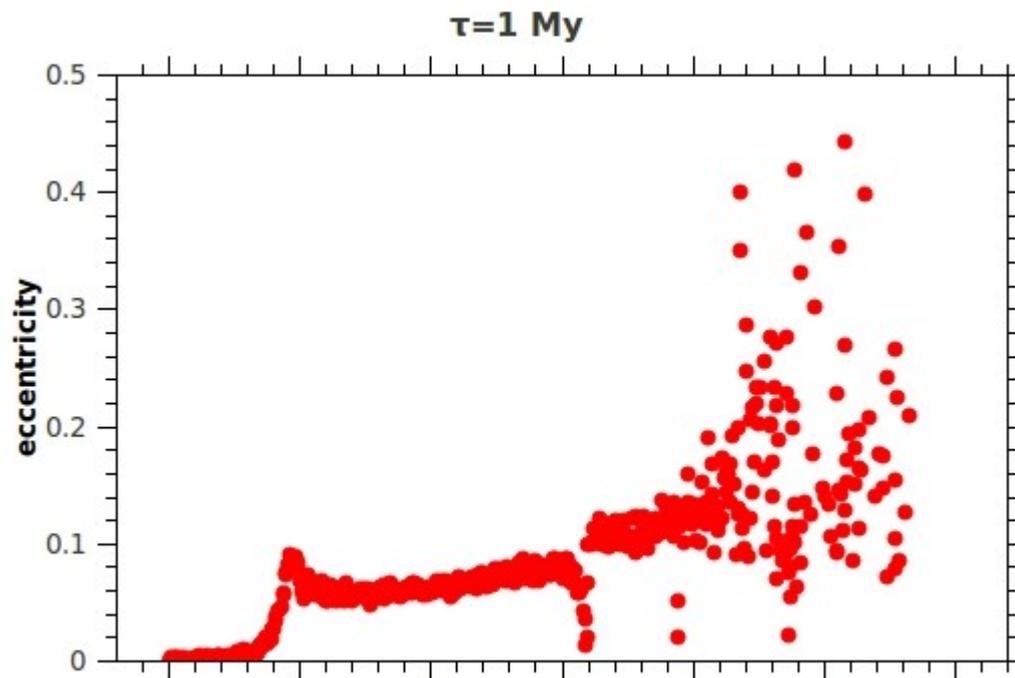
asteroid at 1.83 AU ($\tau=3$ My)



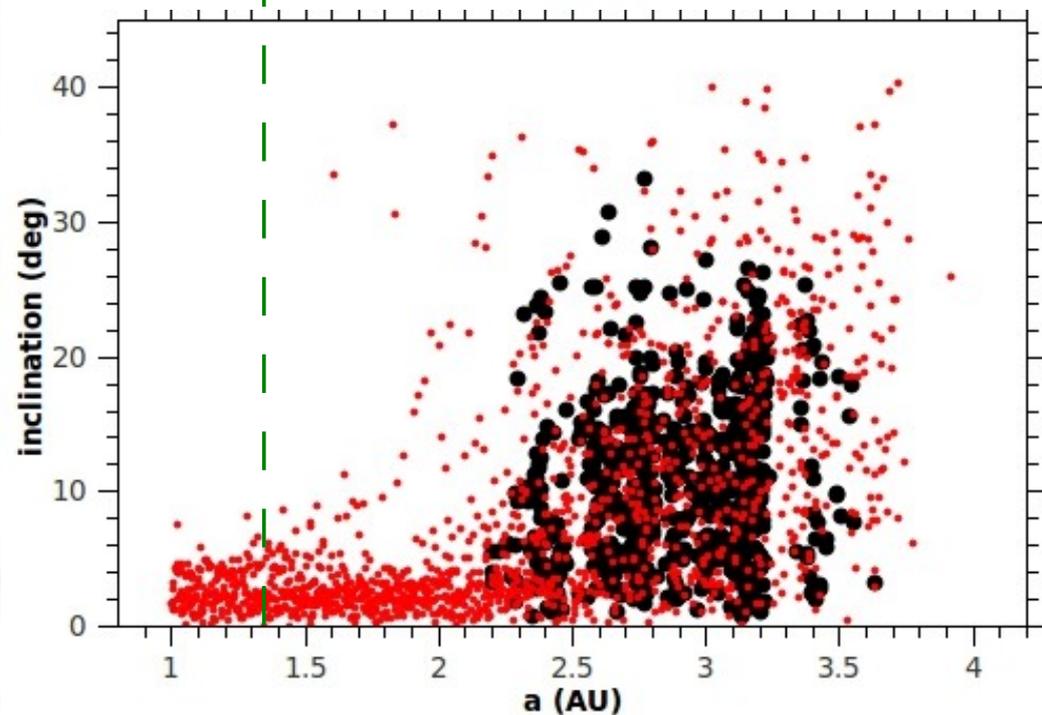
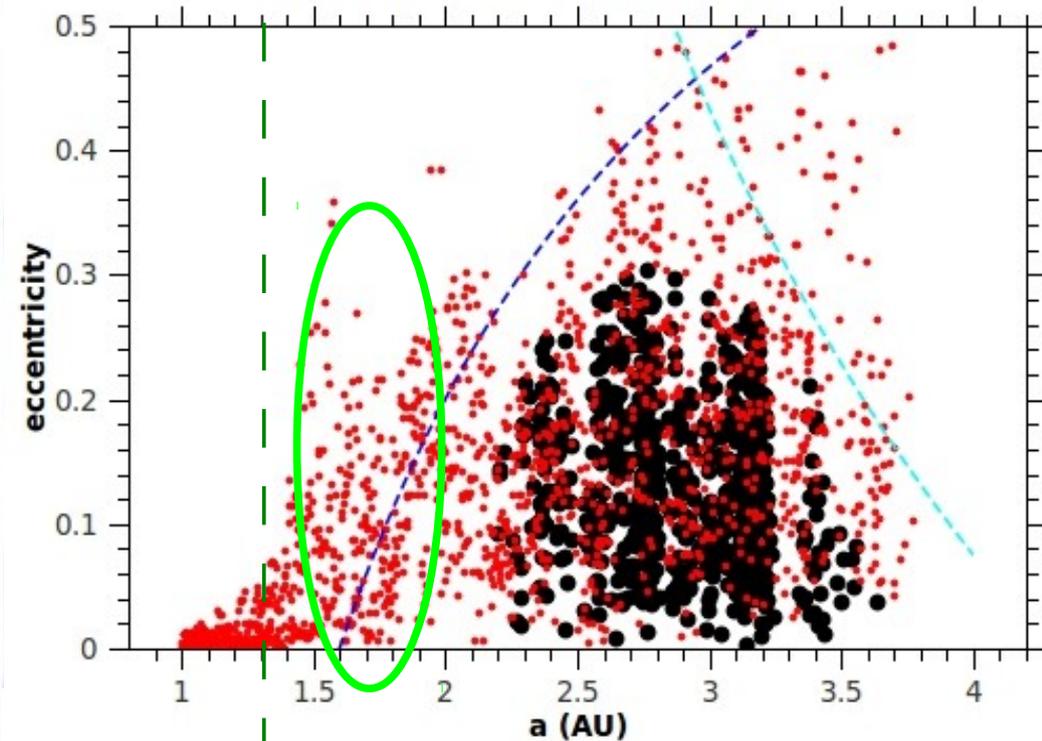
Results (1)

- Several episodes of **pericentric** and **nodal** SRs
 - asteroids can go “**up and down**” in e, i
- Strongest eccentricity forcing when $g \sim 16$ "/yr
- Inclination increases strongly when $e > 0.1$

Results (2)



Real vs. Simulated asteroids



Results (3)

- (a,e) and (a,i) distributions well reproduced
→ no spurious “gaps” or “clusters”
- Critical $\tau \sim 3$ My
- Very little mass loss ($< 9\%$)
- Regions near $a = 2$ AU (v6 SR), outside the “wedge”, or inside the MMRs are unstable
- An excited Extended-belt is created
- The disk forms an Outer edge ~ 1.3 AU → small- q orbits

Conclusions

- The orbital distribution of main-belt asteroids could be the result of the **slow evaporation of the proto-solar nebula**
 - **multiple SR crossings** scramble the eccentricities and inclinations of asteroid orbits
 - these SRs are directly related to the **initial configuration of the planets** (and the disk mass/geometry) **
- An excited, Extended (**E-**)belt is created ($a < 2$ AU) → maybe needed as a **source of LHB projectiles**
- The disk of solids is **truncated** at ~ 1.2 AU → important for explaining **the small mass of Mars** in terrestrial planet formation simulations